



NORTHWESTERN
UNIVERSITY

Superfluid ^3He in Aerogel

W.P. Halperin
Northwestern University

H. Choi, G. Gervais, K. Yawata, T. M. Haard,
Northwestern University

J.A. Sauls, Northwestern Univ.

N. Mulders, Univ. of Delaware

R. Nomura, Tokyo Inst. Tech.

Y. Lee, Univ. of Florida

Moses Chan, Penn. State Univ.

National Science Foundation DMR-0244099

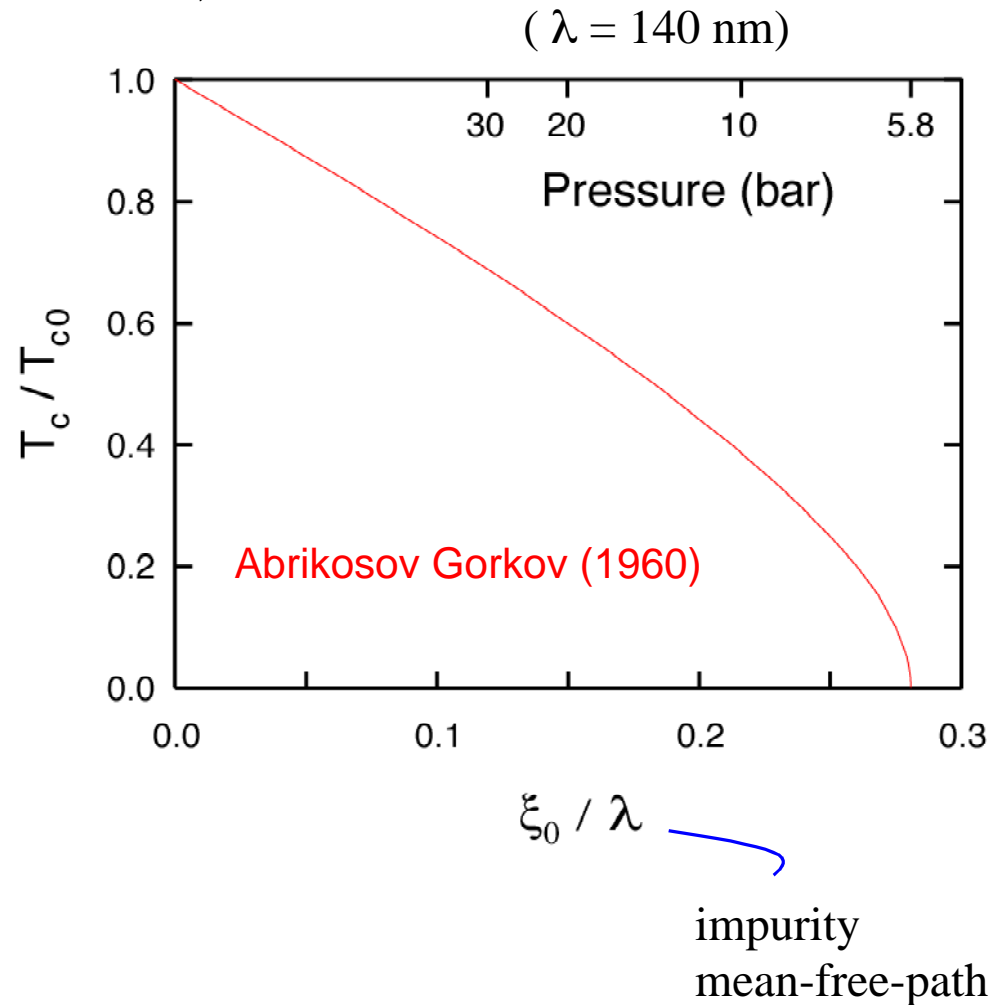
Impurities \Rightarrow Disorder

Unconventional Superfluids :

- heavy fermions (UPt₃), ruthenates,
- cuprates
- ³He

Impurity suppression
of *unconventional*
superfluidity

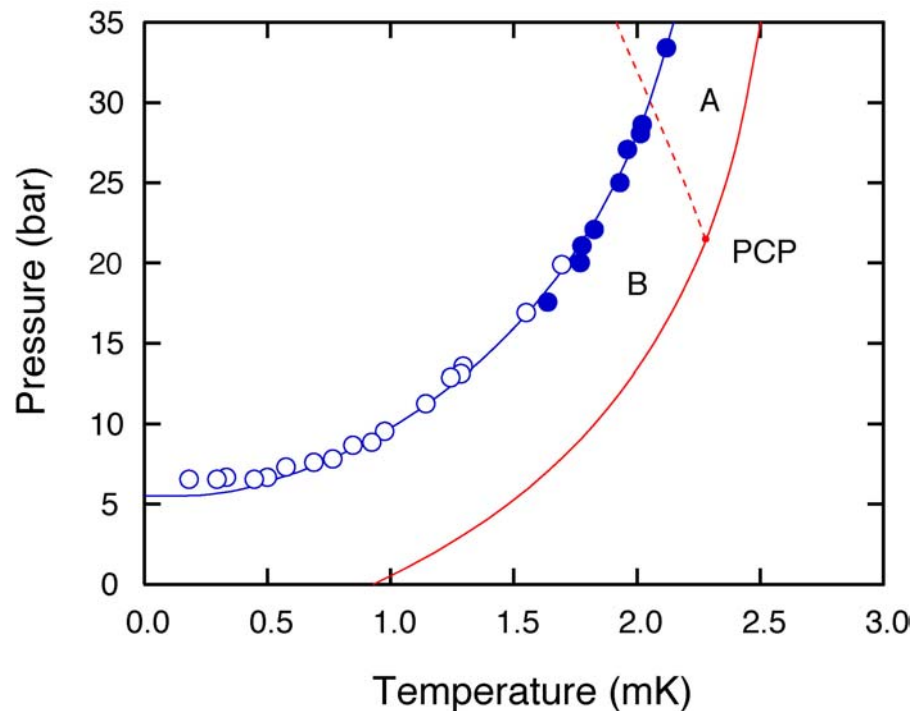
Disorder can be introduced in
superfluid ³He by
high porosity silica aerogel





Why Superfluid ^3He in Aerogel ?

- Superfluid ^3He in aerogel; hopelessly complex, or *perturbatively simple*?
- Can we learn about *pure* superfluid ^3He by introducing spatial disorder?
- Do the aerogel superfluids show *new phenomena*? new phases, metastability, and nucleation.



Matsumoto et al. PRL 79, 253 (1997)
Gervais et al. PRL 87, 035701 (2001);
PRB 66, 054528 (2002).

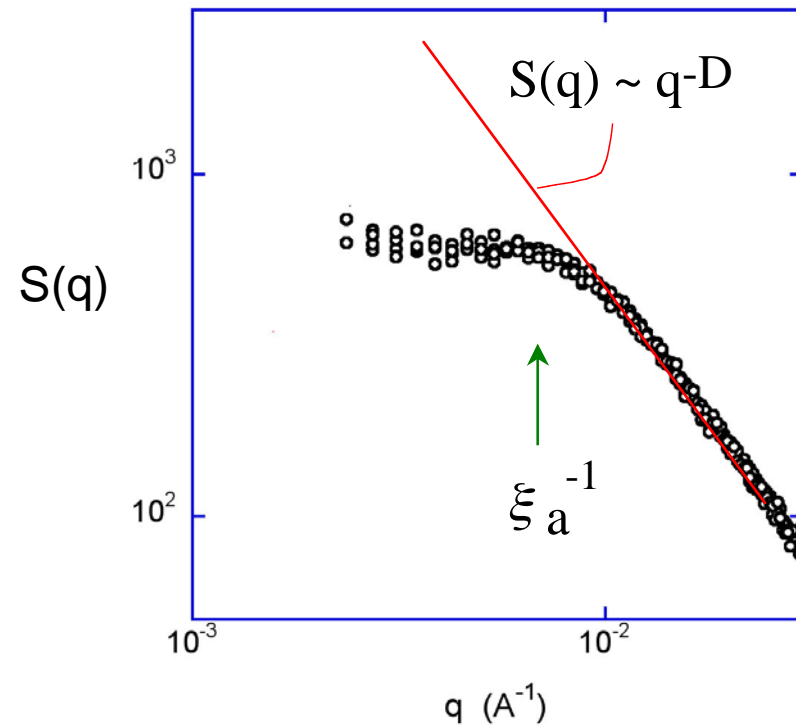
Silica Aerogel:

SAXS,

porosity, $\phi = 98\%$, *Norbert Mulders*

$$S(q) \sim q^{-D}, \quad q > \xi_a^{-1}$$

$$S(0) \sim \xi_a^3 (1 - \phi)^2$$



DLCA Calculation,

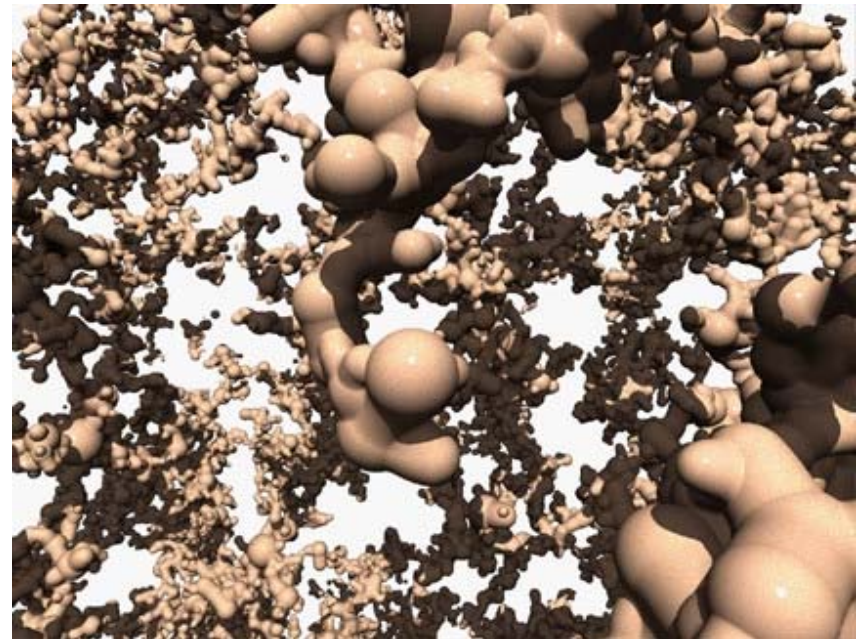
porosity, $\phi = 98\%$, *Tom Haard*

particle diameter $\delta \approx 3$ nm

mean free path $\lambda \approx 200$ nm

Fractal cutoff $\xi_a \approx 30 - 50$ nm

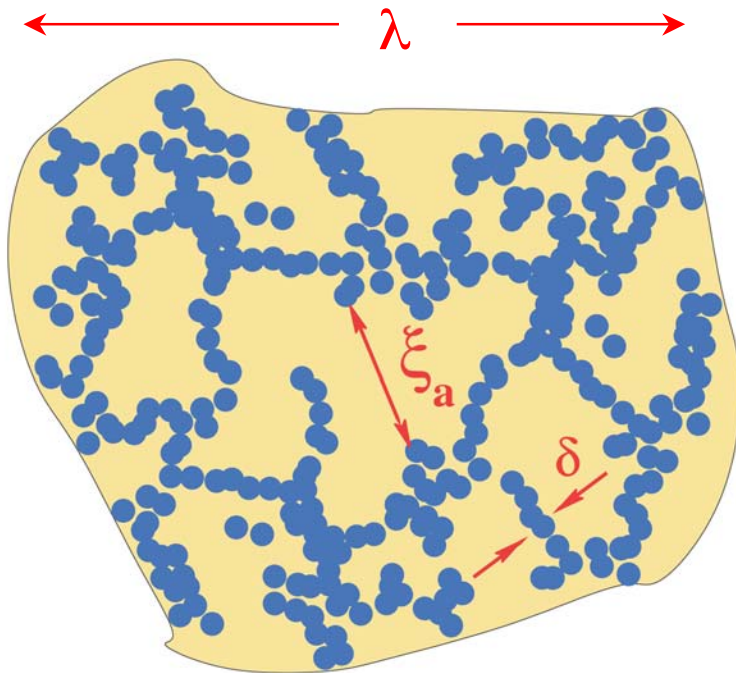
Fractal dimension $D \approx 1.8$



Compare these length scales with the **superfluid coherence length**:

$$\xi_0 \approx 20 \text{ to } 80 \text{ nm (34 to 0 bar)}$$

The coherence length is larger than the silica particle size diameter, δ ;
it is of order the correlation length, ξ_a ;
and smaller than the mean free path, λ (superfluid is not completely suppressed)



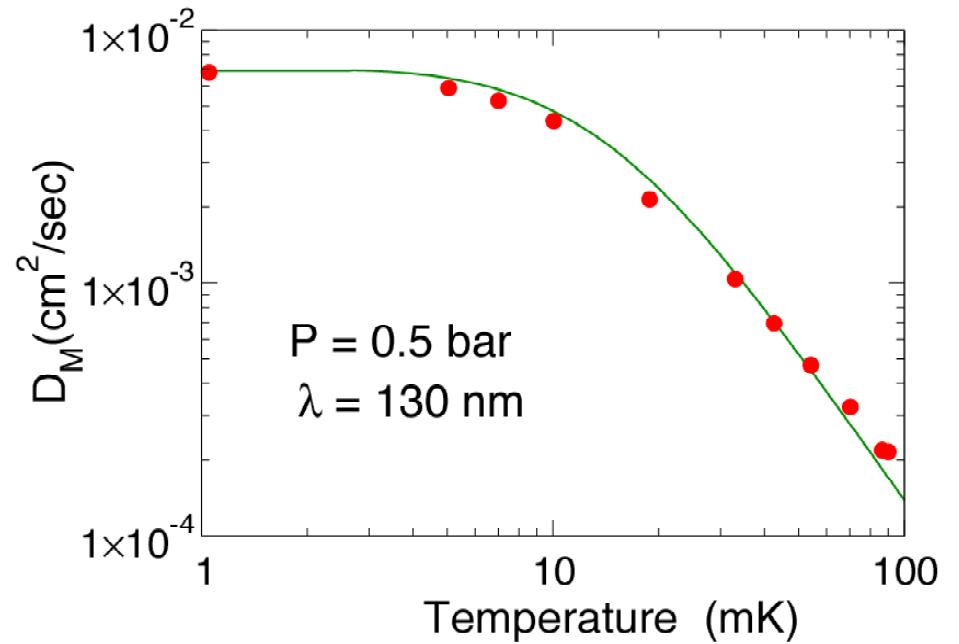
The mean free path can be determined in the **normal state** of ^3He in aerogel from transport measurements:

- spin diffusion (U. Mass., Grenoble)
- thermal conductivity (Lancaster)
- attenuation of sound (Northwestern)

Normal State

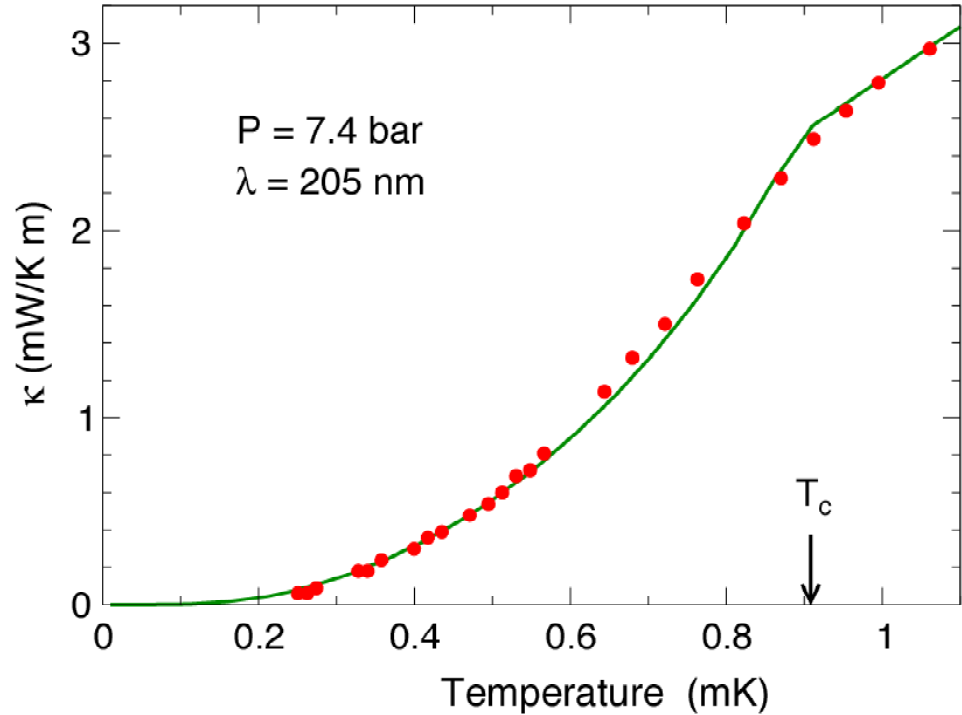
Spin diffusion

Collin, thesis (2002);
Sauls *et al.* cond-mat/0408492



Thermal conductivity

Fisher *et al.* JLTP (2001)



Scattering Models for the Superfluid State:

Thuneberg et al. PRL 80, 2861 (1998); Sauls, Sharma, PRB (2003); Hanninen, Thuneberg PRB (2003).

The simplest model for aerogel: a random distribution of point isotropic scatterers.

Arbitrarily small anisotropy leads to orientational disorder of the vector order parameter, according to fundamental arguments by Imry and Ma, PRL(1975). For the axial state (bulk A-phase), this forces spatial inhomogeneity of the orbital state which Volovik JETPL (1997) calls a superfluid glass.

Fomin JETPL (2003); JLTP (2004) has suggested that instead of the axial state an aerogel A-phase, could be a *homogeneous* equal spin pairing state, but one of a class of “robust” *p*-wave states not constrained by above arguments to be a glass.

The free energy, **density of states**, thermal conductivity, attenuation of sound, spin diffusion, magnetic susceptibility have been calculated. Note: *Thuneberg, Sauls, Khardze, Nagai, Mineev, Fomin.*

Ginzburg-Landau theory:

For *pure* ^3He the free energy $F[A]$ is expanded as a function of the five fourth-order invariants of the 3×3 complex p -wave order parameter, A :

$$F[A] = \frac{1}{3} \alpha(T) \text{Tr}(AA^\dagger) + g_z B(AA^\dagger)B + \\ \beta_1 |\text{Tr}(AA^T)|^2 + \beta_2 (\text{Tr}(AA^T))^2 + \beta_3 \text{Tr}[AA^T(AA^T)^*] \\ \beta_4 \text{Tr}[(AA^\dagger)^2] + \beta_5 \text{Tr}[AA^\dagger(AA^\dagger)^*]$$

Rainer and Serene PRB **13**, 4745 (1976); **17**, 2901 (1978).

where $\beta_{ijk} = \beta_i + \beta_j + \beta_k$

axial state (A-phase) heat capacity jump:

$$\frac{\Delta C_A}{C} = 1.19 \frac{2}{\beta_{245}}$$

isotropic state (B-phase) heat capacity jump:

$$\frac{\Delta C_B}{C} = 1.43 \frac{5}{3\beta_{12} + \beta_{345}}$$

A_1 - A_2 field splitting ratio:

$$r = \frac{T_{A1} - T_c}{T_c - T_{A2}} = -\frac{\beta_5}{\beta_{245}}$$

T_{AB} suppression by magnetic field:

$$1 - \frac{T_{AB}(B)}{T_c} = g(\beta) (B/B_0)^2$$

$$g(\beta) = \frac{\beta_{245}}{2(-3\beta_{13} + 2\beta_{345})} \left(1 + \sqrt{\frac{(3\beta_{12} + \beta_{345})(2\beta_{13} - \beta_{345})}{\beta_{245}\beta_{345}}} \right)$$

GL-theory for Superfluid ^3He in Aerogel

Homogeneous Isotropic Scattering Model (HISM):

Treat the silica aerogel as a system of point scatterers for ^3He quasiparticles. There is one parameter: the transport mean free path, λ , assuming strong scattering (unitary limit). The pair breaking parameter is $x = \xi_0 / \lambda$.

Homogeneous Isotropic Scattering Model (HISM) *Thuneberg et al. PRL 80, 2861 (1998)*; *Sauls (2002)*, includes **strong coupling** corrections:

$$\beta_i^{aero} = a_1(T_{ca}, \lambda) \beta_i^{wc} + \frac{T_{ca}}{T_{c0}} \beta_i^{sc-bulk}$$

Inhomogeneous Isotropic Scattering Model (IISM):

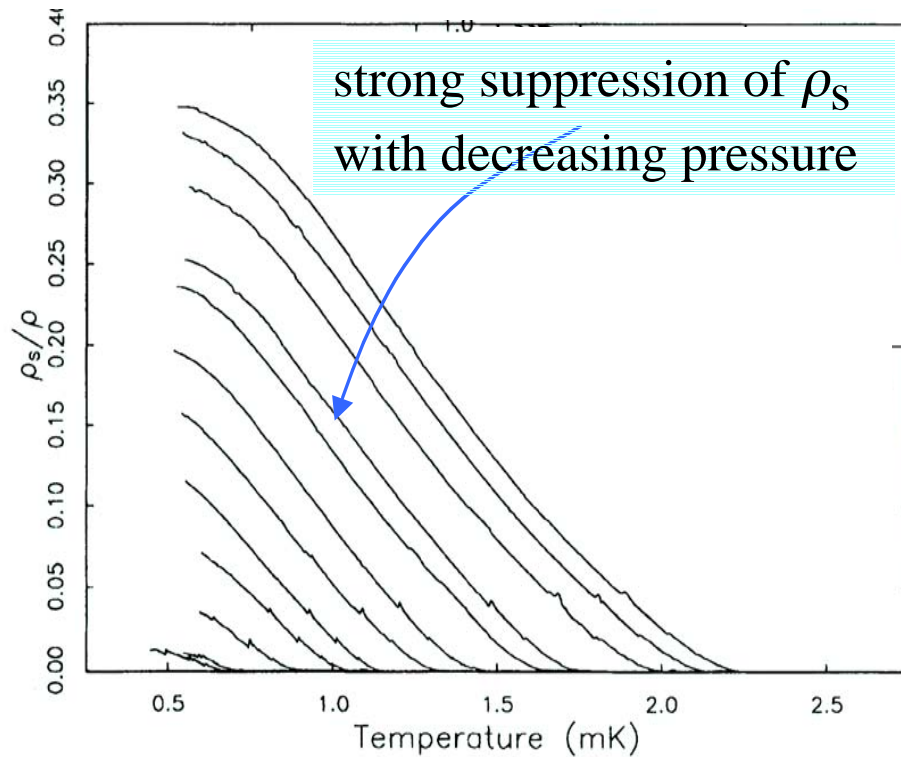
According to Sauls and Sharma (2003), include the particle-particle correlation length, ξ_a , in a new pair breaking parameter, $x = (\xi_0 / \lambda) / [1 + (\xi_a / \xi_0)^2]$.

$$\beta_i^{aero} = a_1(\lambda, \xi_a) \beta_i^{wc} + \frac{T_{ca}}{T_{c0}} \beta_i^{sc-bulk}$$

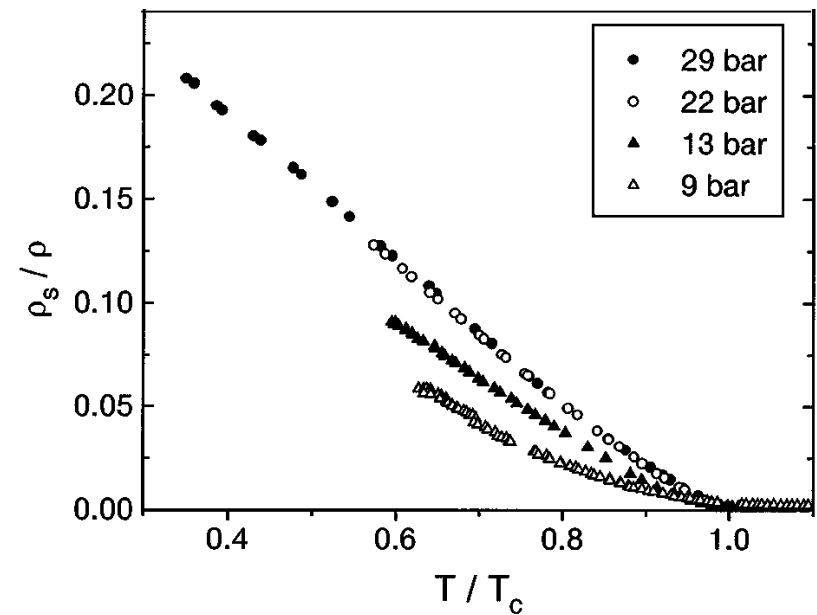
Experiments for Superfluid ^3He in Aerogel:

Superfluid density *Porto and Parpia PRL 74, 4667 (1995). Golov et al. PRL 82, 3492 (1999).*

torsional oscillator

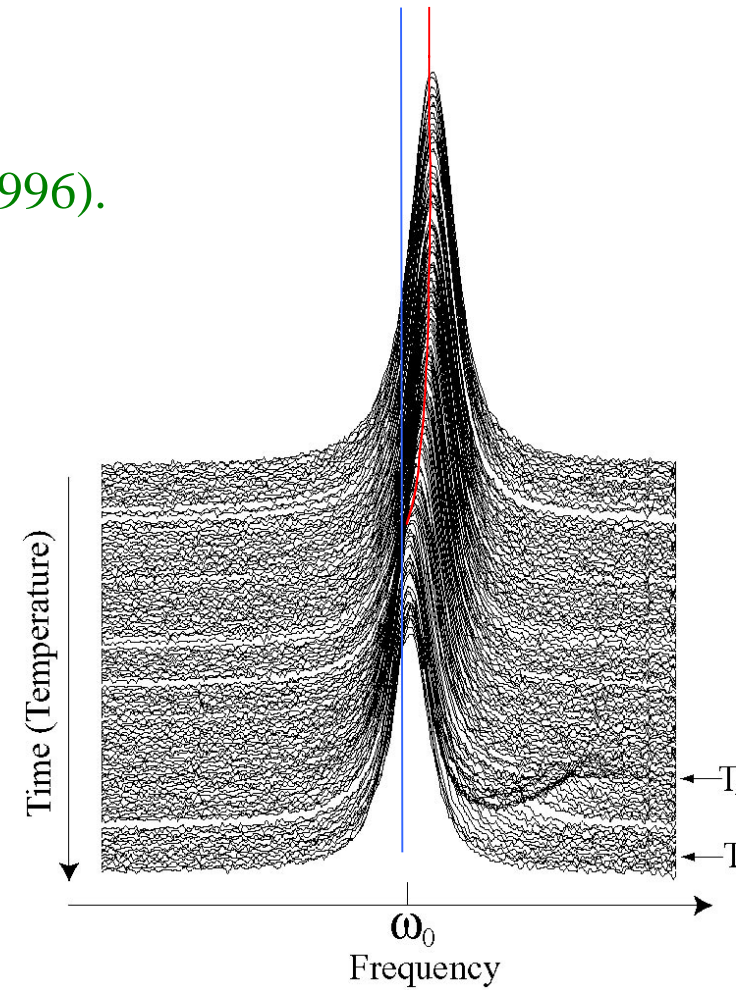
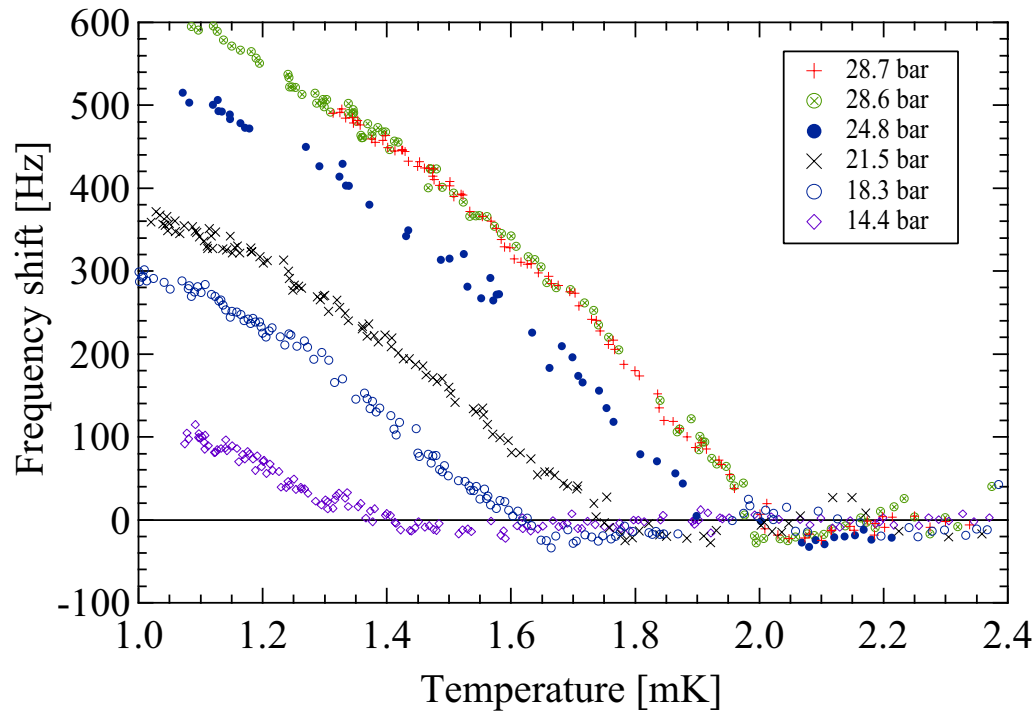


sound velocity



NMR

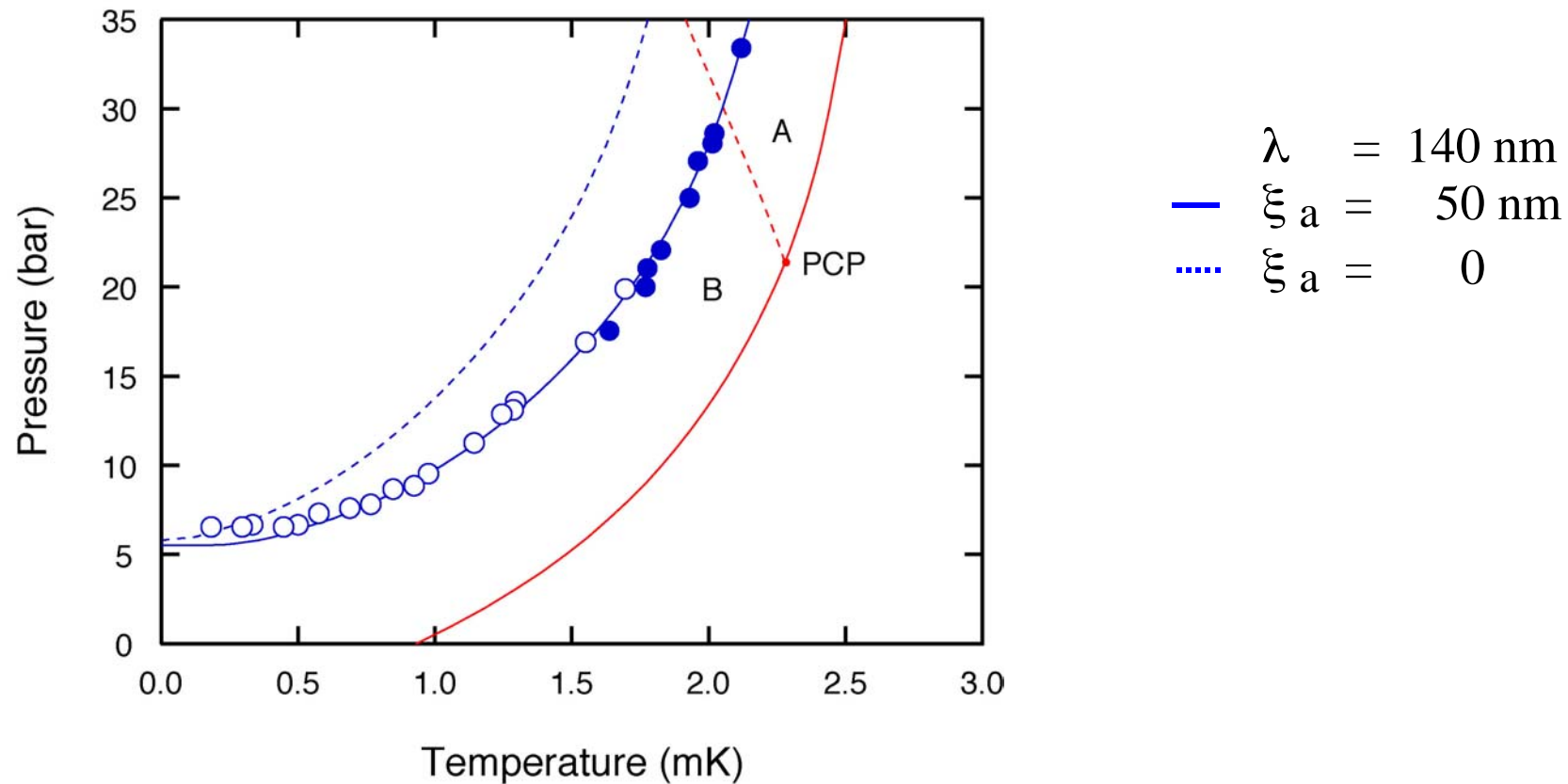
Sprague et al. PRL 75, 661 (1995); PRL 77, 4568 (1996).



P = 28.7 bar

More recently: *Osheroff (Stanford), Hata (Osaka), Bunkov/Godfrin, (Grenoble), Dimitriev (Moscow)*

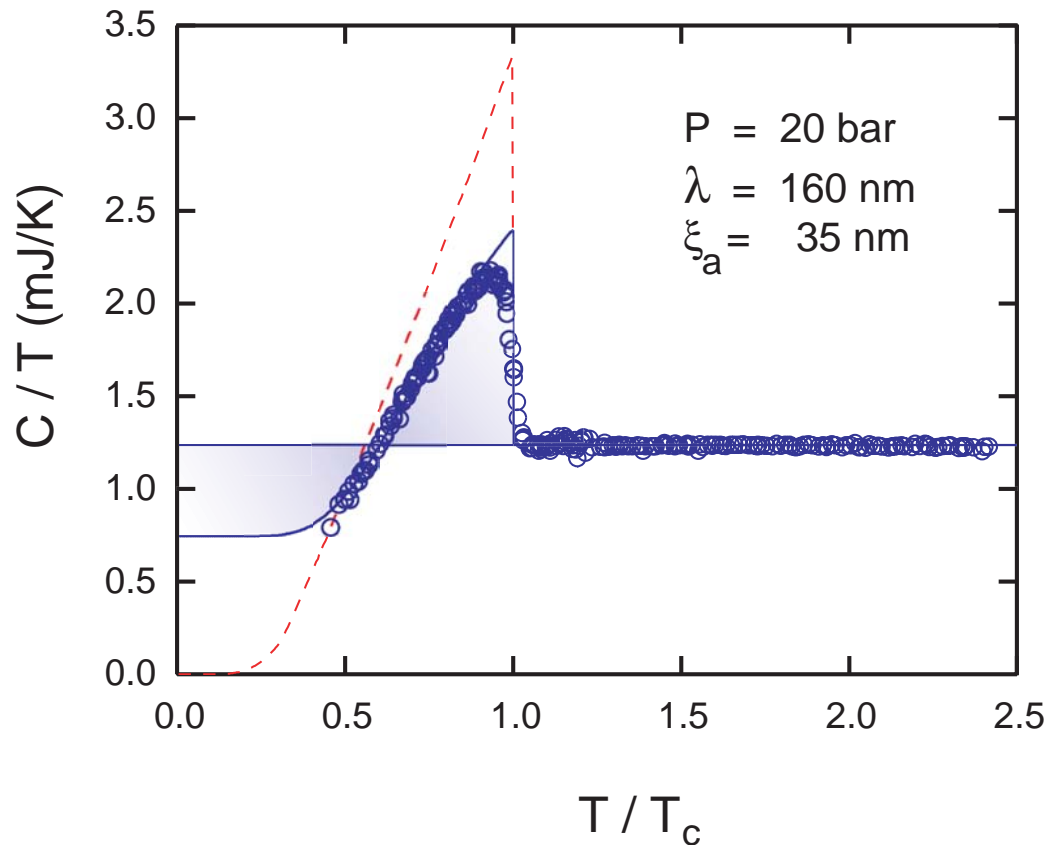
Phase Diagram with correlated scattering (IISM)



- ➔ Is this a well-defined thermodynamic transition?
- ➔ Is the phase boundary field dependent?
- ➔ What are the phases of this aerogel/superfluid: axial (A), isotropic (B), planar, polar, axi-planar,....?

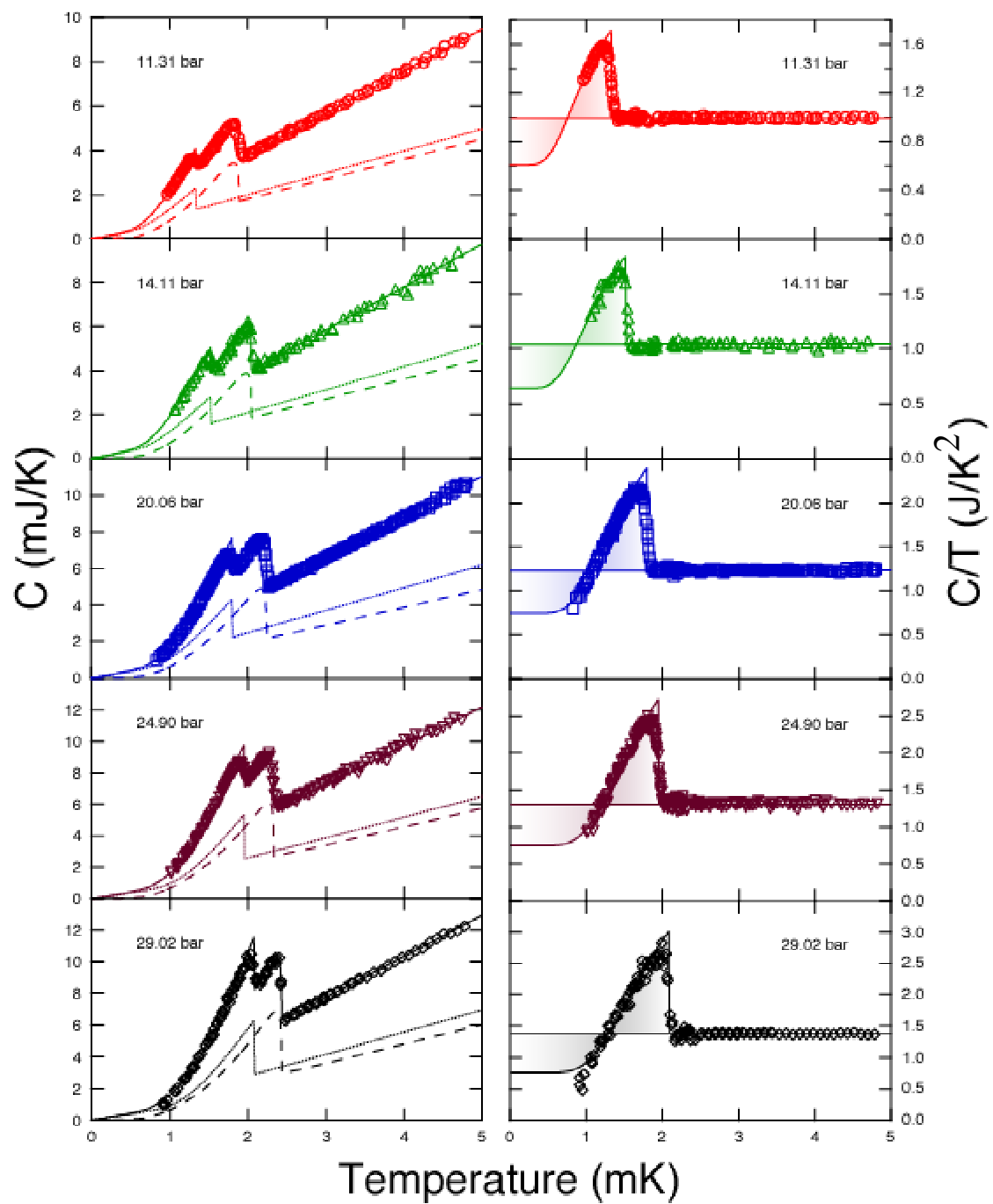
Heat Capacity:

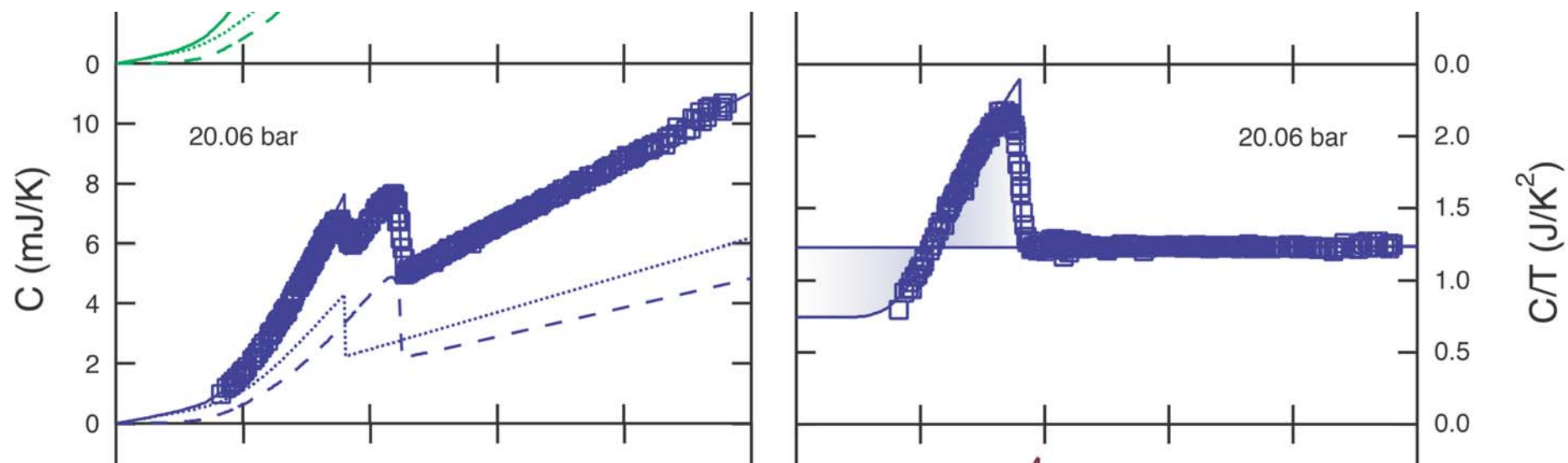
Aerogel-Superfluid heat capacity at $P = 20$ bar
(98% porosity aerogel) Choi *et al.* PRL **93**, 145301 (2004).



$$\Delta(T) = \pi k_B T_c \left[\frac{2}{3} \left(\frac{T_c}{T} - 1 \right) \frac{\Delta C}{C} \right]^{1/2}$$

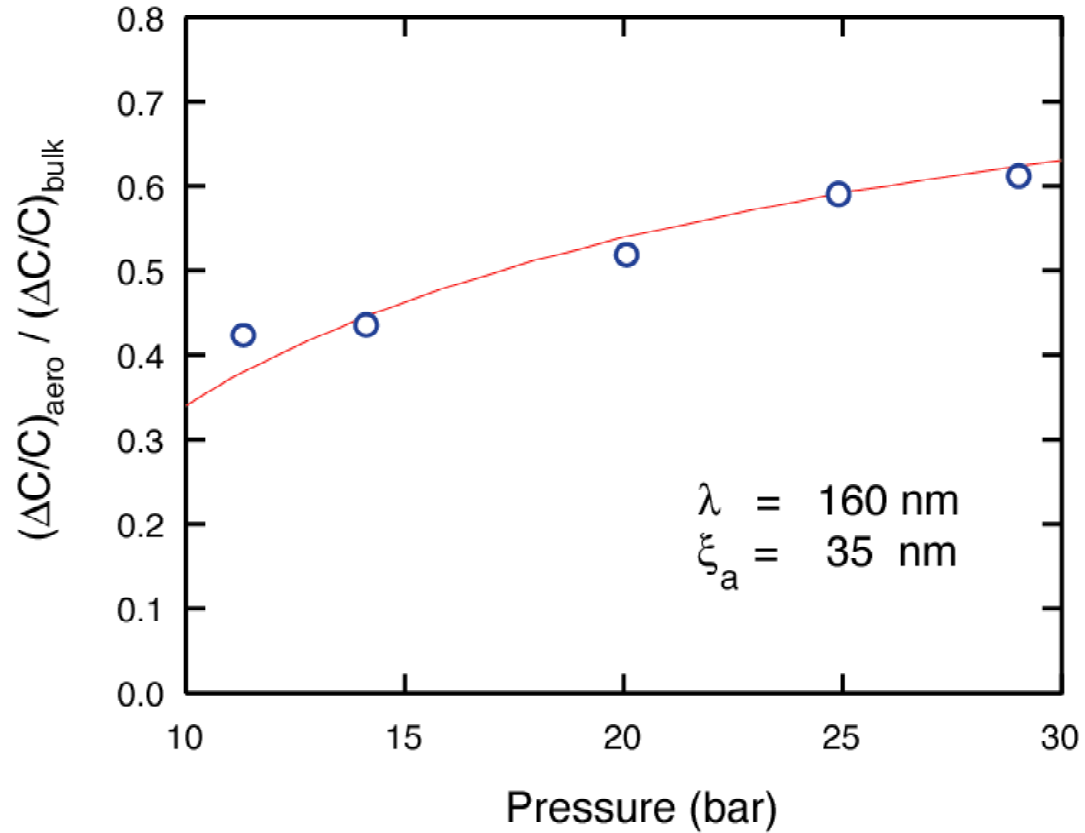
$\Delta C/C = 1.0 \pm 0.1$ compared to 1.82 for pure ^3He at $P = 20$ bar.
the order parameter is suppressed.





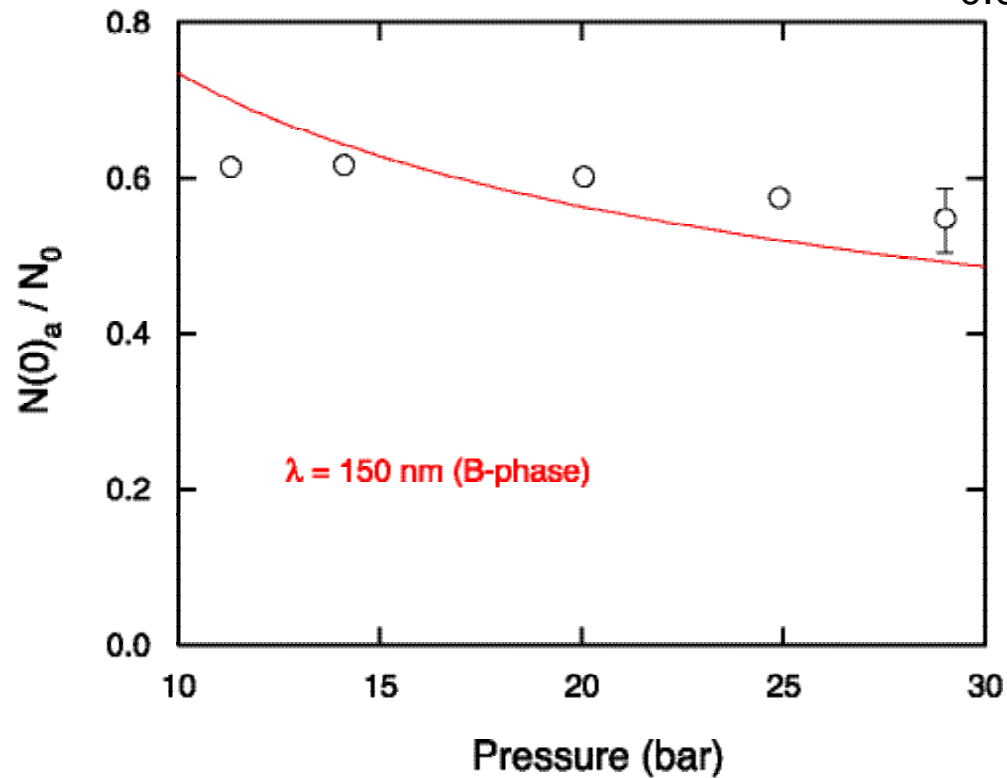
Strong suppression of the amplitude of the order parameter:

$$\Delta(T) = \pi k_B T_c \left[\frac{2(T_c - T)}{3T} \frac{\Delta C}{C} \right]^{1/2}$$

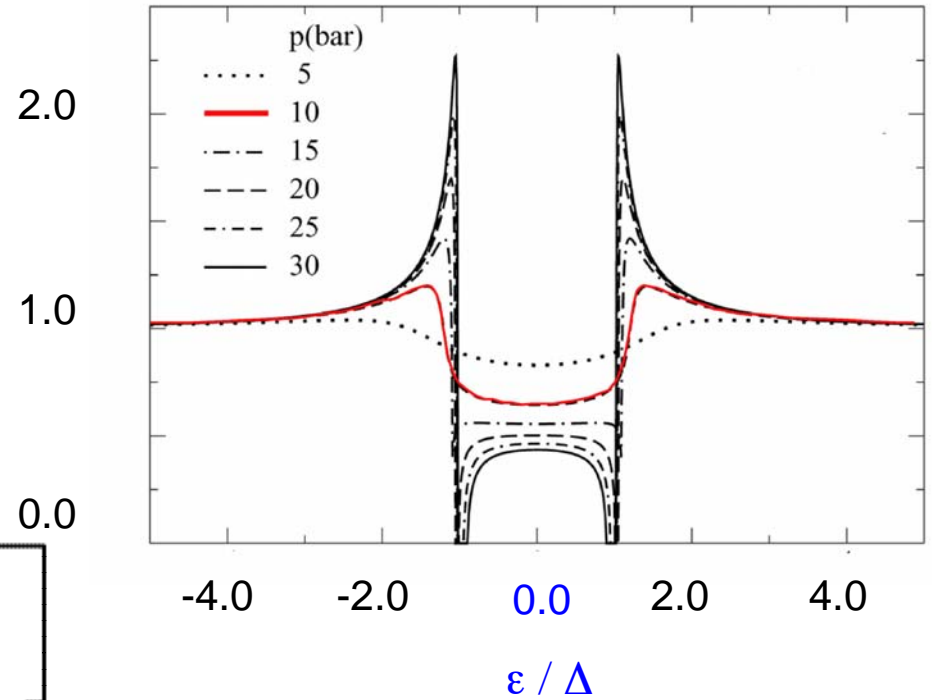


Density of Andreev bound states
at $\epsilon = 0$

Sharma and Sauls
JLTP (2001).



$N_a(\epsilon) / N_0$

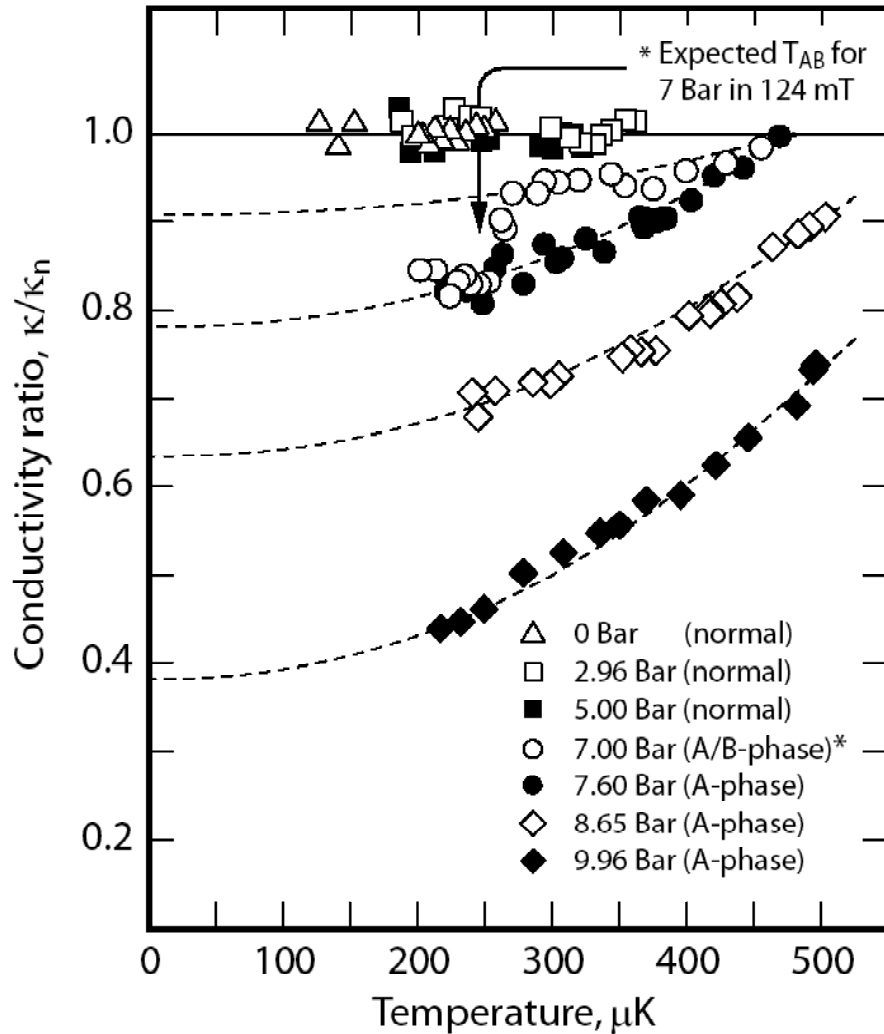


From heat capacity,
Choi et al. PRL (2004).

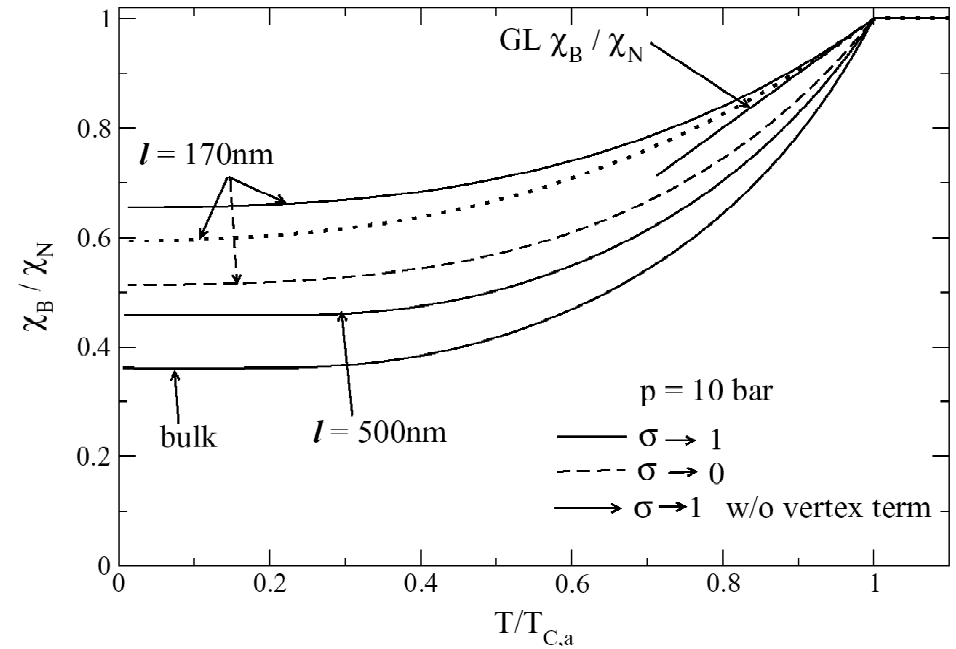
Disorder induces gapless
superfluidity

Magnetic Susceptibility

Sharma and Sauls, JLTP **125**, 115 (2001).

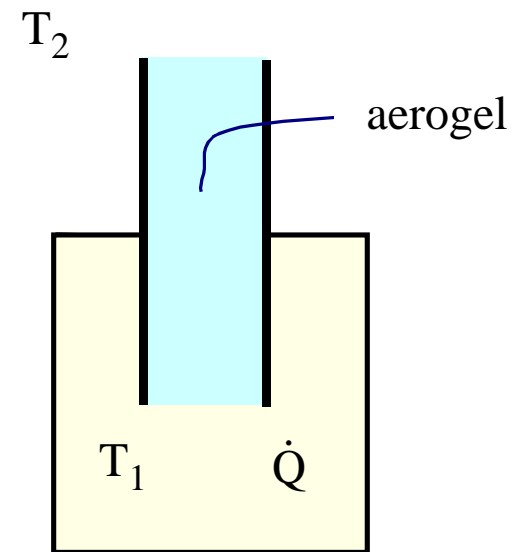


CosLab 2004; Chamrousse, December 18, 2004



Thermal Conductivity

Fisher et al. PRL **91**, 105303 (2003)

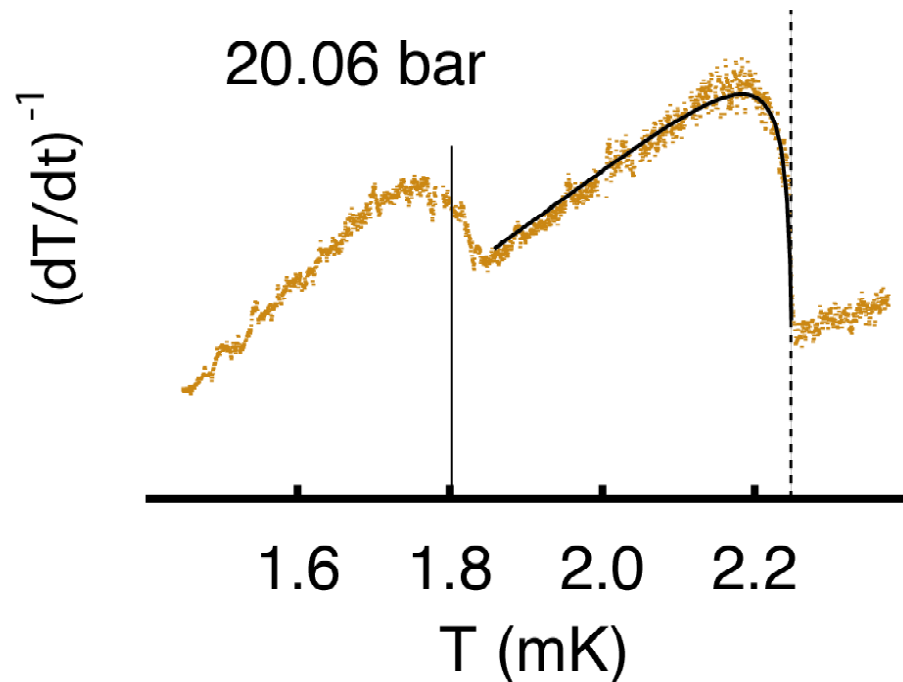
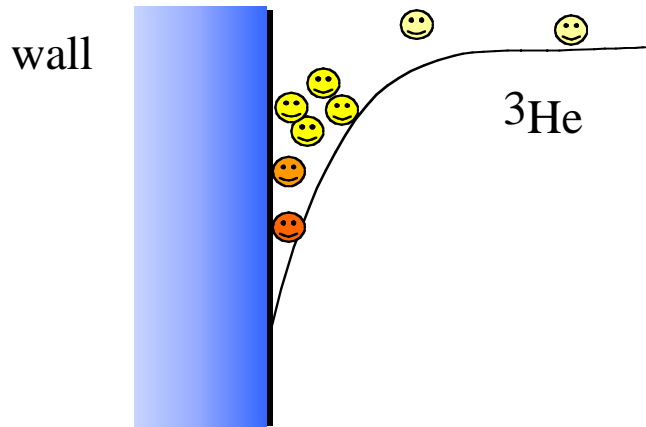


Andreev surface bound states from bulk specific heat:

Andreev bound states have been predicted at surfaces in ^3He by many workers including: Kjalman, Kukijarvi, Rainer; Buchholz, Rainer; Zhang; Vorontsov and Sauls;...

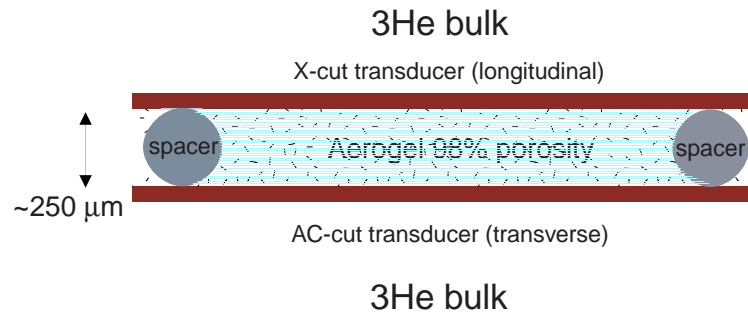
We have measured the heat capacity of Andreev *surface* bound states at the bulk transition temperature (dashed line) showing rounding from the temperature dependence of the coherence length:

Slow warm-up specific heat:



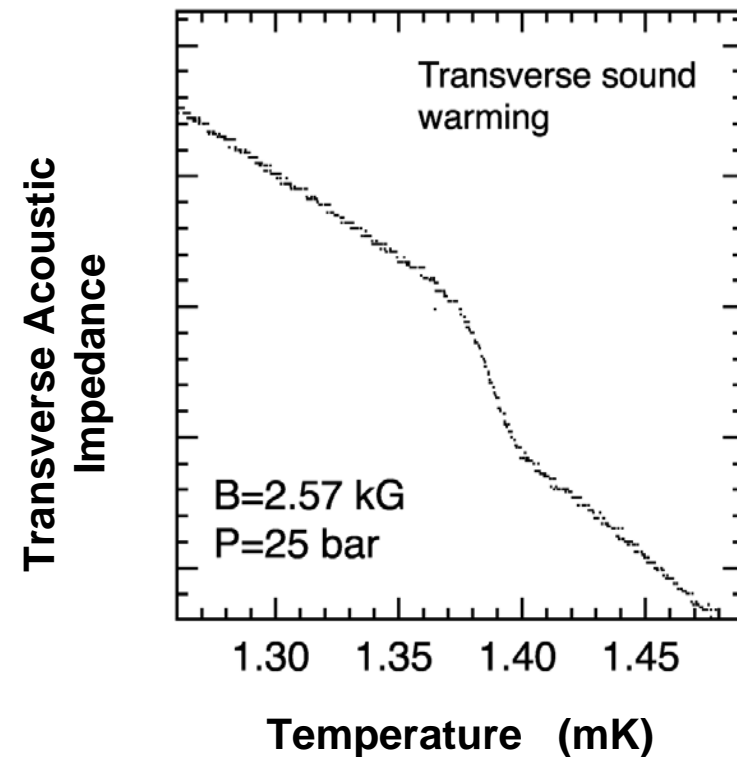
Phase Diagram:

Aerogel/Superfluid ^3He phase diagram from transverse acoustic impedance:



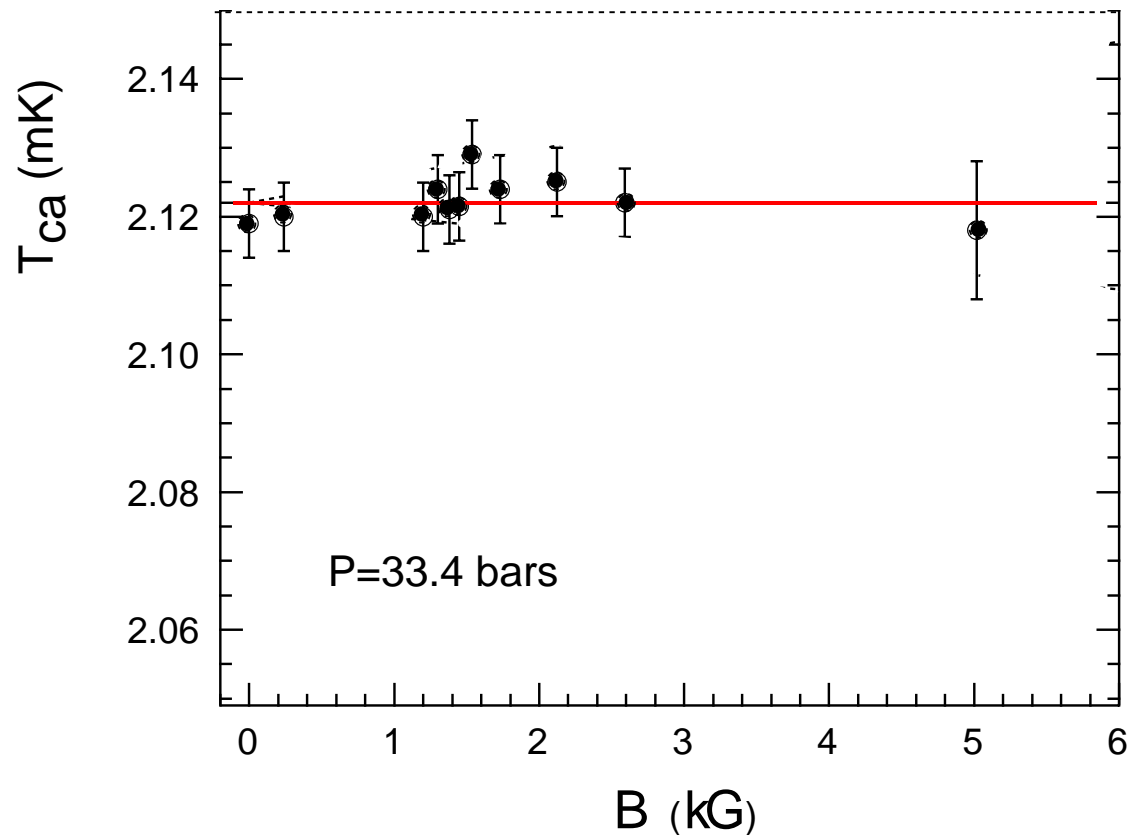
B- to A-transition in a magnetic field

We measure the acoustic impedance and find a response at all known phase transitions: bulk and aerogel superfluids. Measurement from zero applied magnetic field up to 0.8 T.
Gervais et al. PRL **87**, 035701 (2001);
PRB **66**, 054528 (2002).

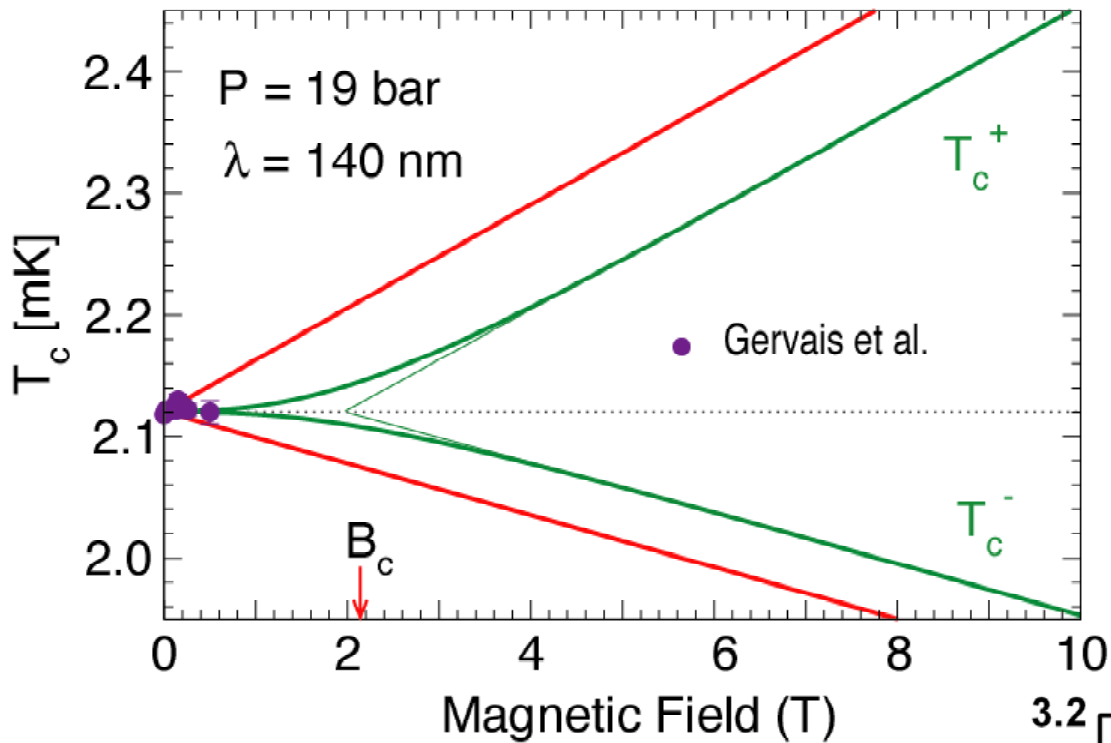


Normal to superfluid transition in a magnetic field:

No field dependence to the aerogel/superfluid transition observed by *Gervais et al. PRB 66,054528 (2002)* up to 0.5 T and no splitting up to 0.8 T



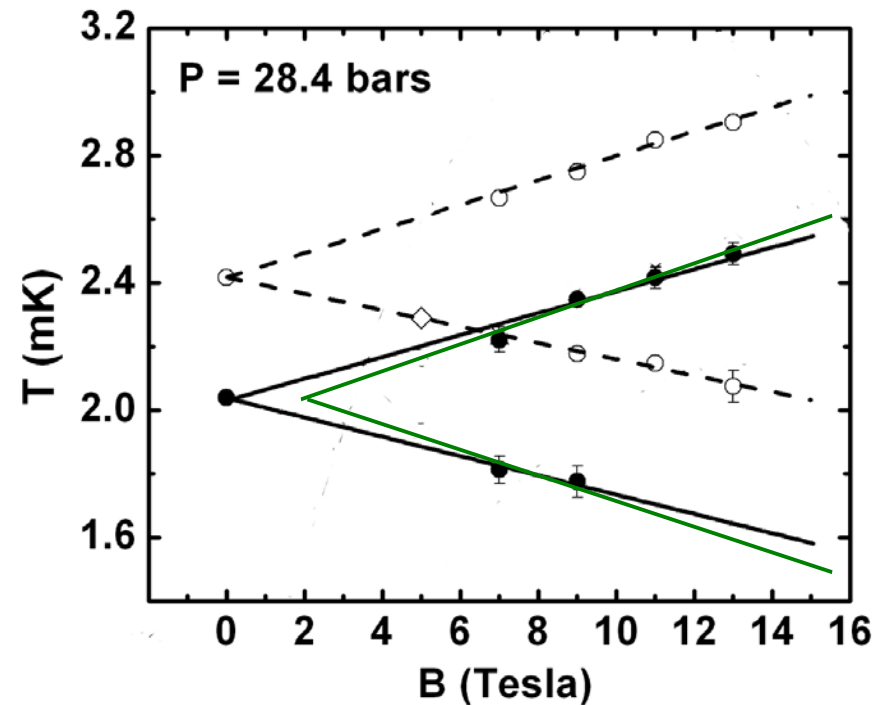
Indicates an equal spin pairing state (ESP).



Antiferromagnetic exchange between liquid ^3He and solid ^3He adsorbed on aerogel give the green curve; *Sauls and Sharma, PRB 68, 224502 (2003)*. Also Baramidze and Kharadze, cond-mat/0307612

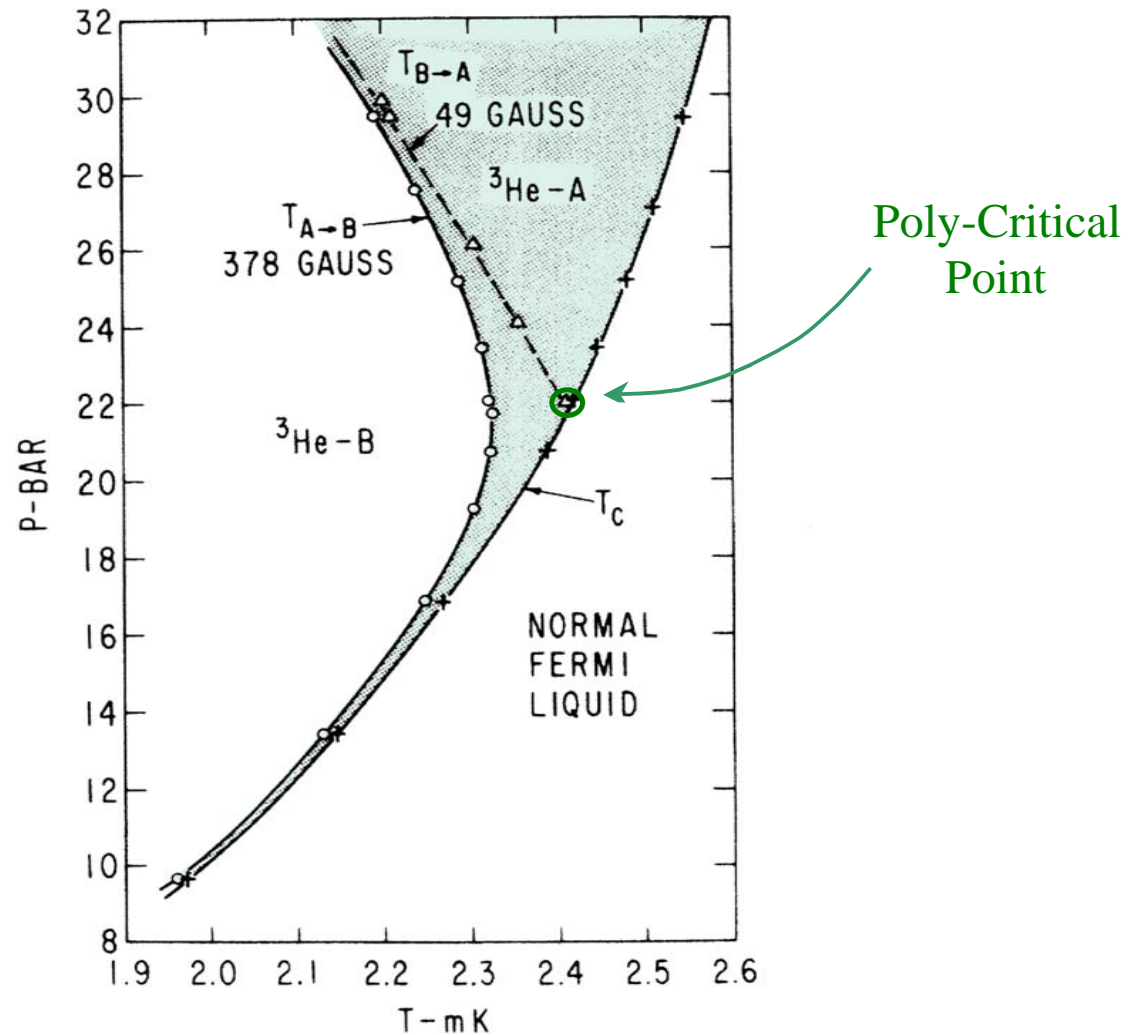
Splitting for aerogel-superfluid discovered at high field:

Choi et al. (Univ. Florida) PRL 93, 145302 (2004)



*AB-transition:
pure superfluid ^3He :*

Paulson et al. PRL 32, 1098 (1974).

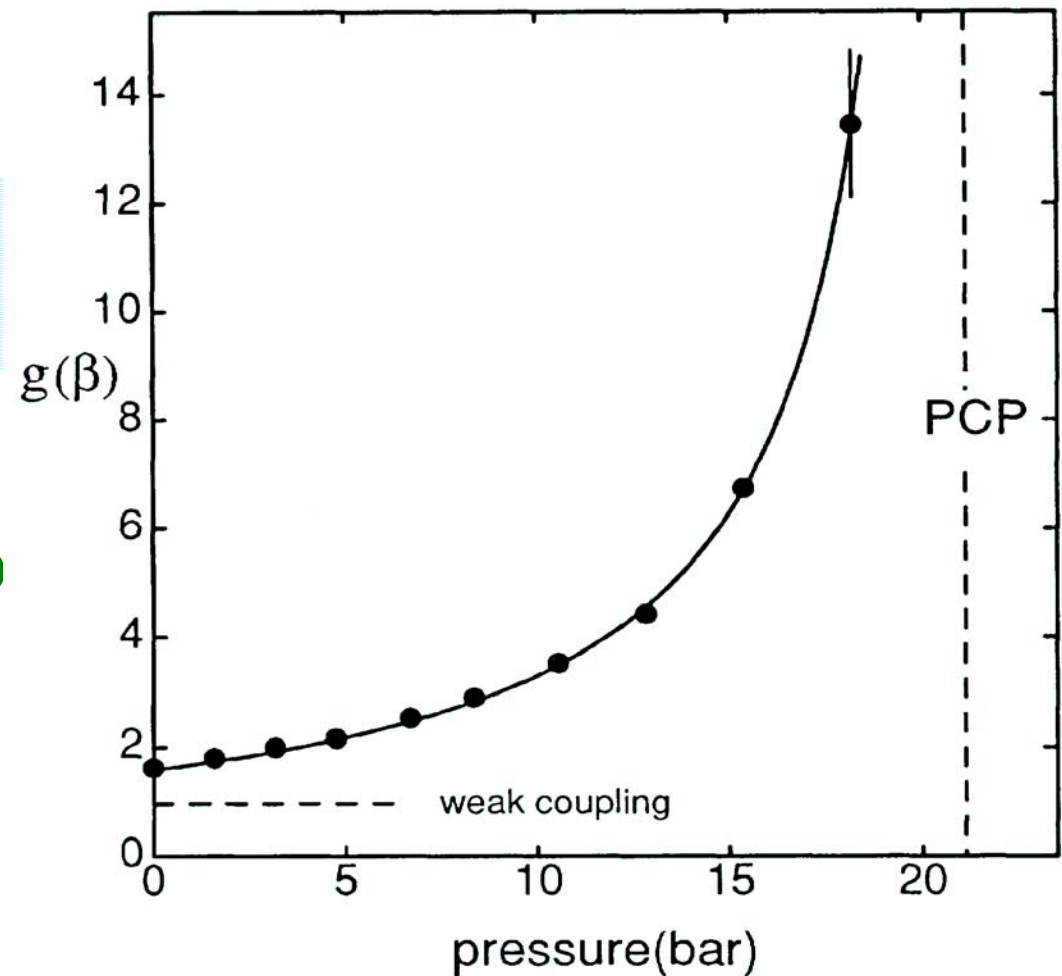


Strong suppression of the B-phase by a magnetic field

⇒ spin triplet state.

$$1 - \frac{T_{AB}(B)}{T_c} = g(\beta) (B/B_0)^2$$

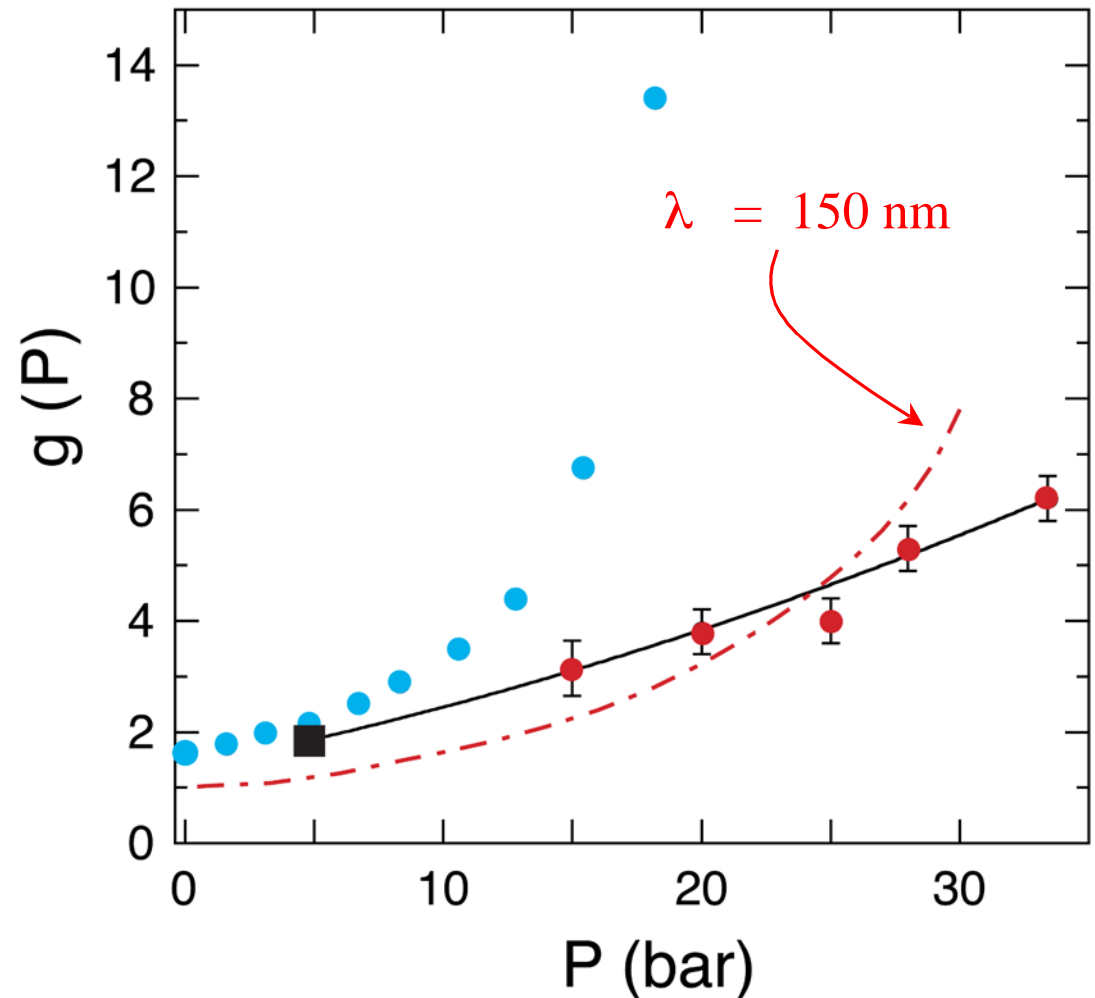
Later detailed study: *Tang, Hahn, Bozler, Gould PRL 67, 1775 (1991)* was performed to determine strong coupling effects: ($\beta_1, \beta_2, \dots, \beta_5, g_z$)



The HISM including strong coupling is in agreement with a vanishing PCP.

- *Gervais et al. PRL 87, 035701 (2001); PRB 66, 054528 (2002)*
- *Brussaard et al. PRL 86, 4580 (2001)*

The B-phase is stabilized by disorder and strong coupling effects are reduced by T_{ca} suppression.

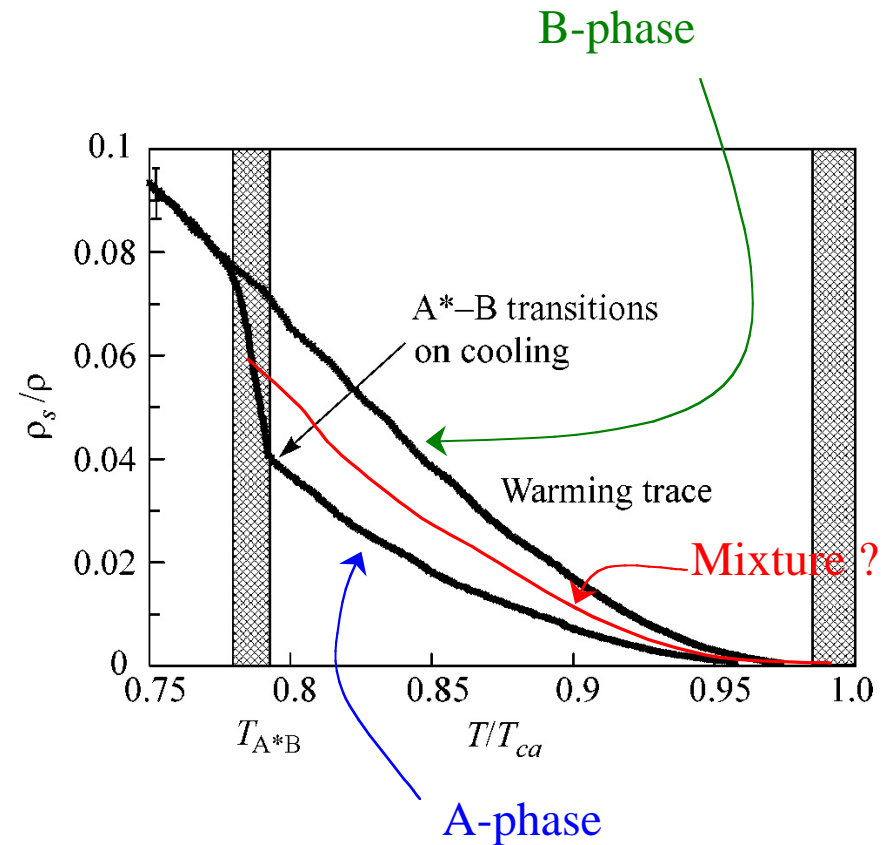
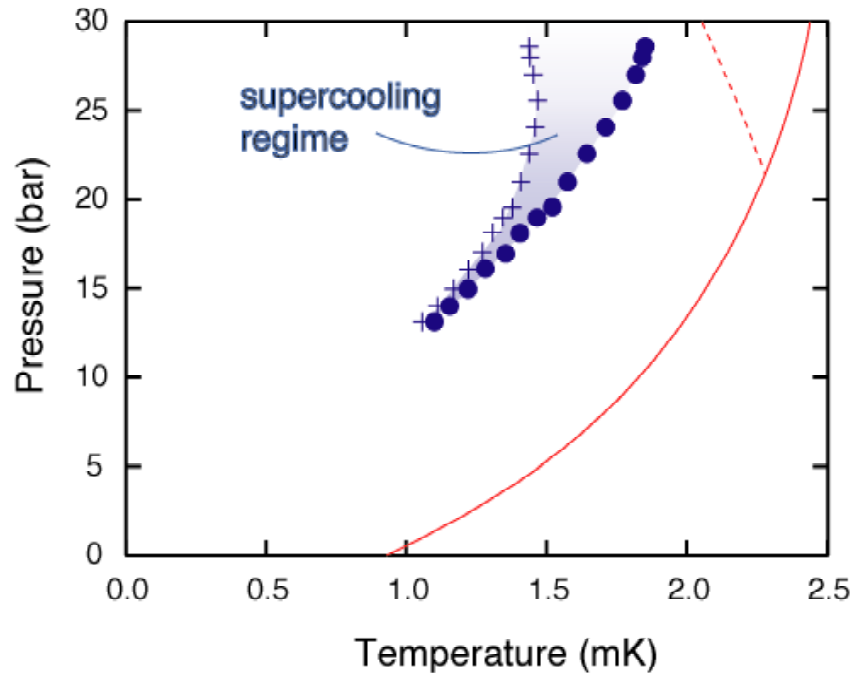


Metastability and Supercooling:

Barker et al., PRL 85, 2148 (2000); Gervais et al., PRL 87, 035701 (2001); PRB 66, 054528 (2002); Baumgardner et al., PRL 85, 2148 (2000); PRL 93, 155301 (2004).

Nazaretski et al., JETP Lett. 79 383 (2004).

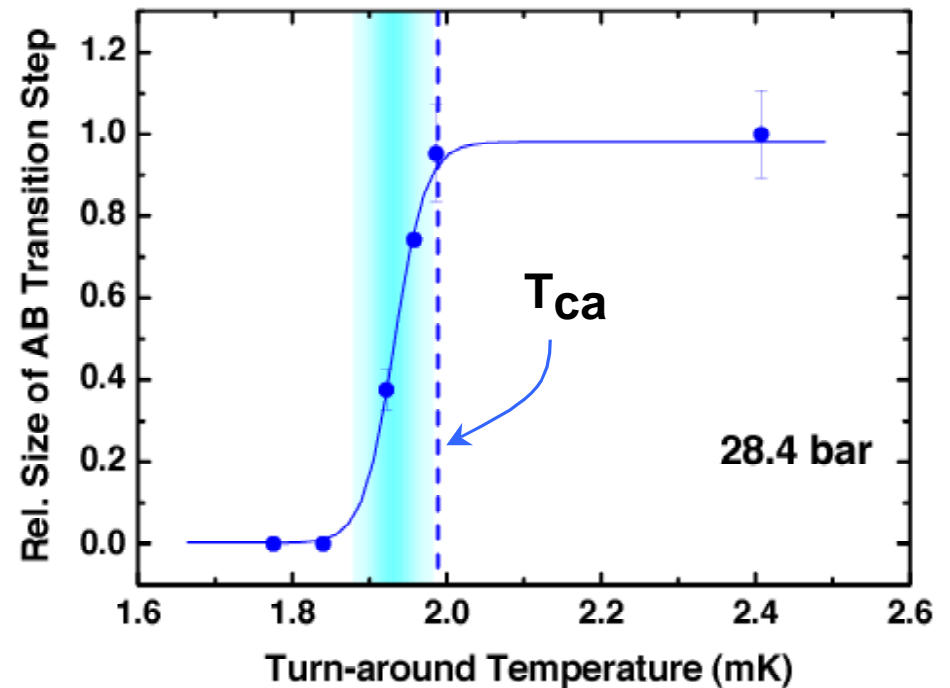
98 % aerogel



The role of anisotropic scattering:

Vicente et al. (Univ. of Florida), submitted.
From transverse acoustic impedance.

There are new terms in the Ginzburg-Landau free energy owing to broken orbital symmetry from anisotropic scattering. *Rainer, Vuorio J.Phys.C (1977); Thuneberg et al., cond-mat/980204; Sauls (2004); Fomin (2004).*



Magnetic field suppression by \mathbf{B} : $f_z = g_z B_\mu A_{\mu i} A_{\nu j}^* B_\nu$
 Anisotropic scattering direction \mathbf{a} : $f_a = \alpha_1 a_i A_{\mu i} A_{\mu j}^* a_j$
 scale is given by an effective field: $B_e = (\alpha_1 / g_z)^{1/2} \approx 0.1 \text{ T}$

PCP is pushed to unobservably high pressure, scattering anisotropy (locally inhomogeneous) will stabilize the A-phase, *Y. Lee (Univ. Florida).*

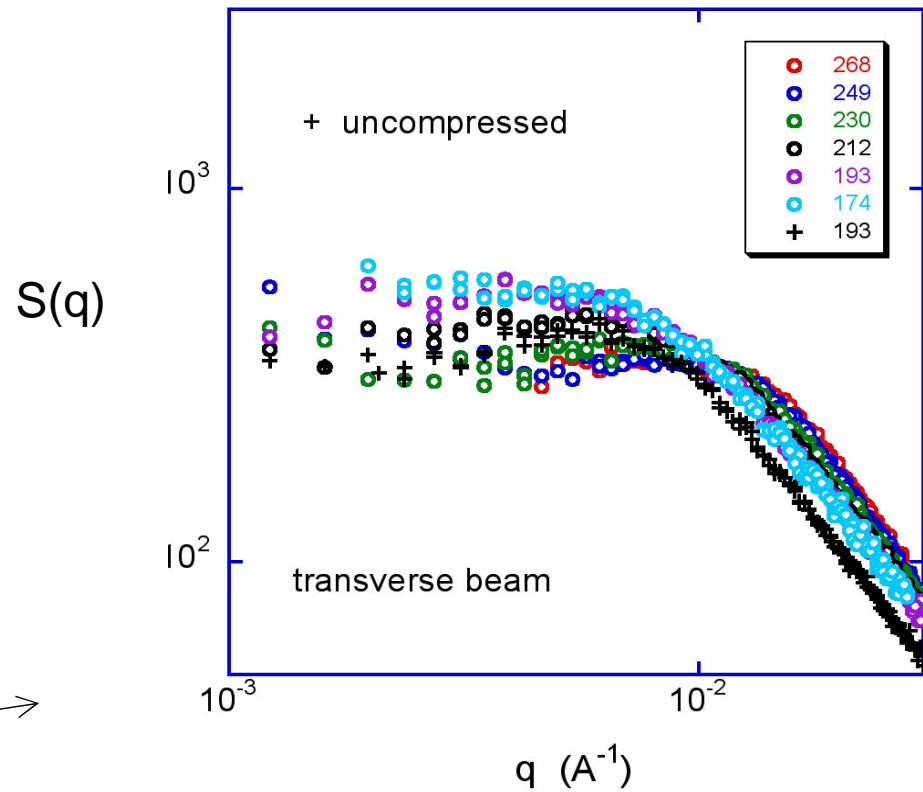
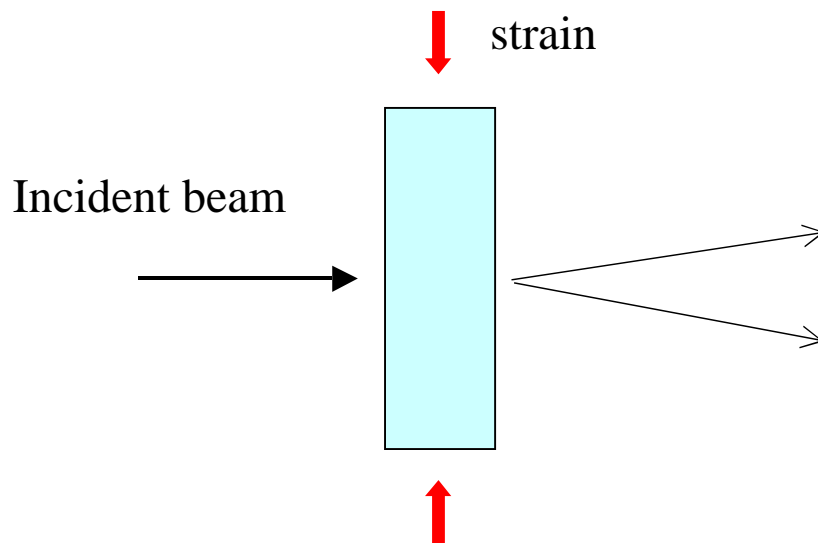
Controlled homogeneous anisotropic scattering:

Introduce uniaxial strain into the aerogel to stabilize the A-phase:

98% aerogel, 30 % strain:

SAXS **Norbert Mulders (2004)**

SAXS:

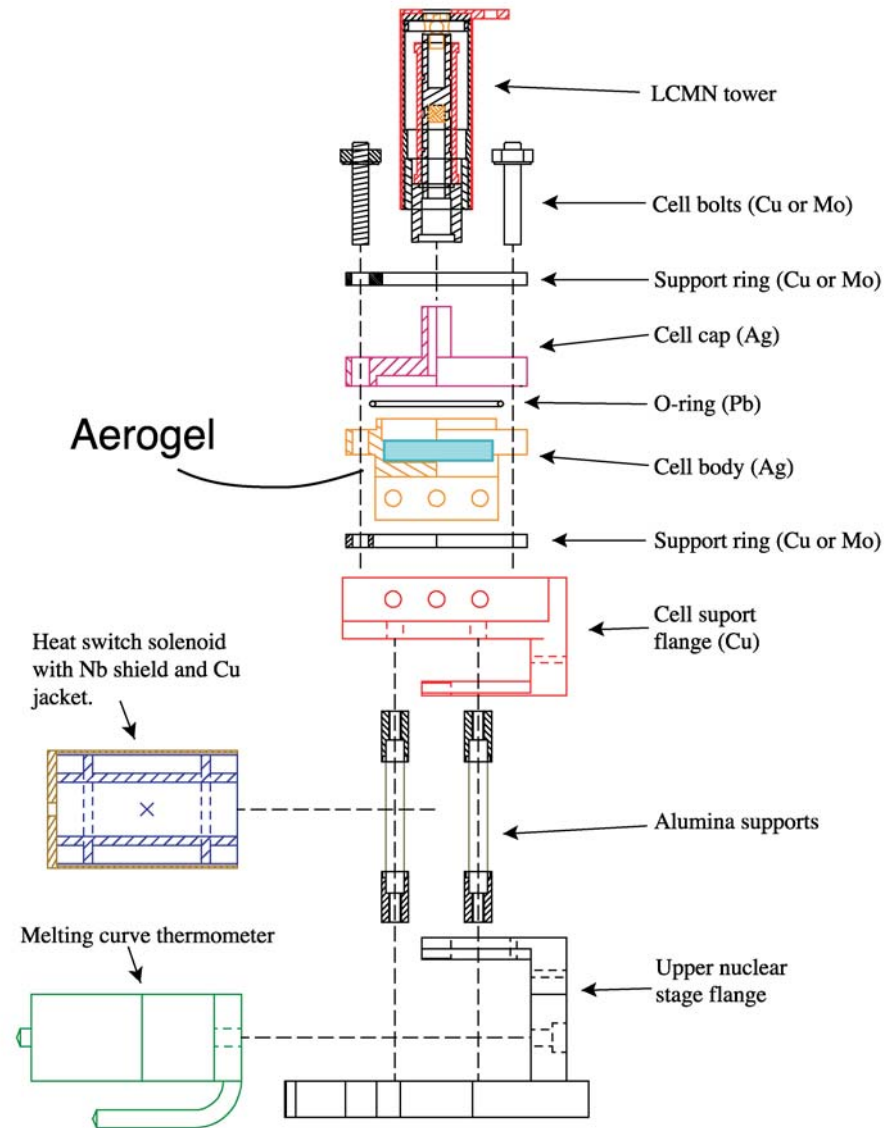


Summary:

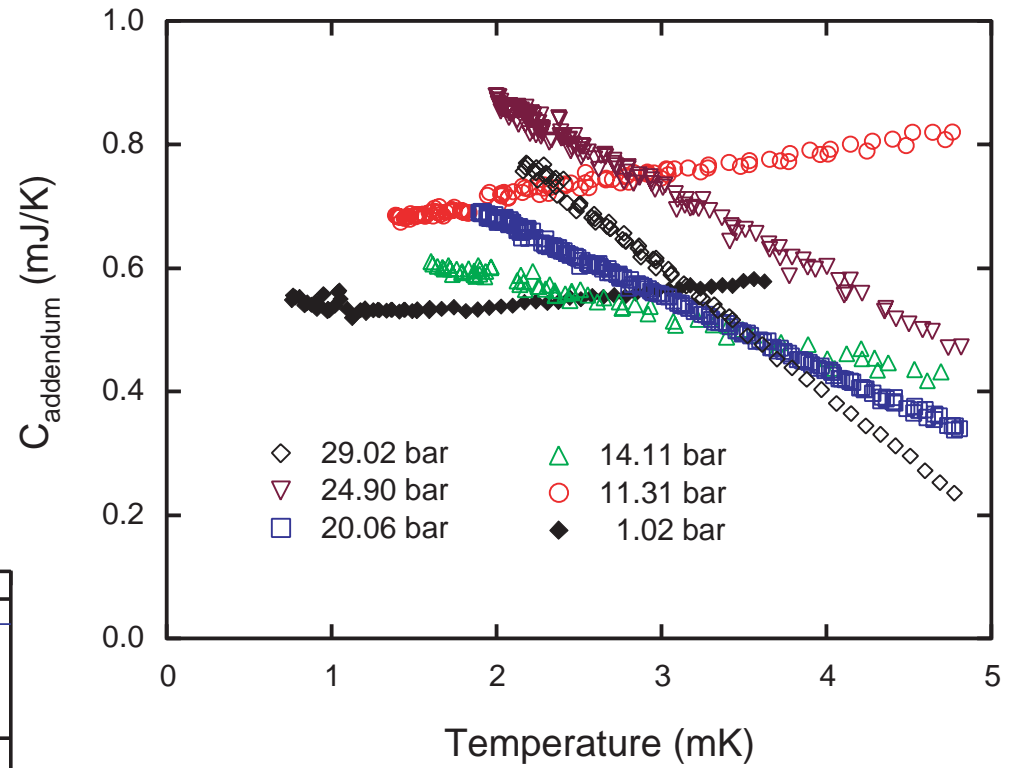
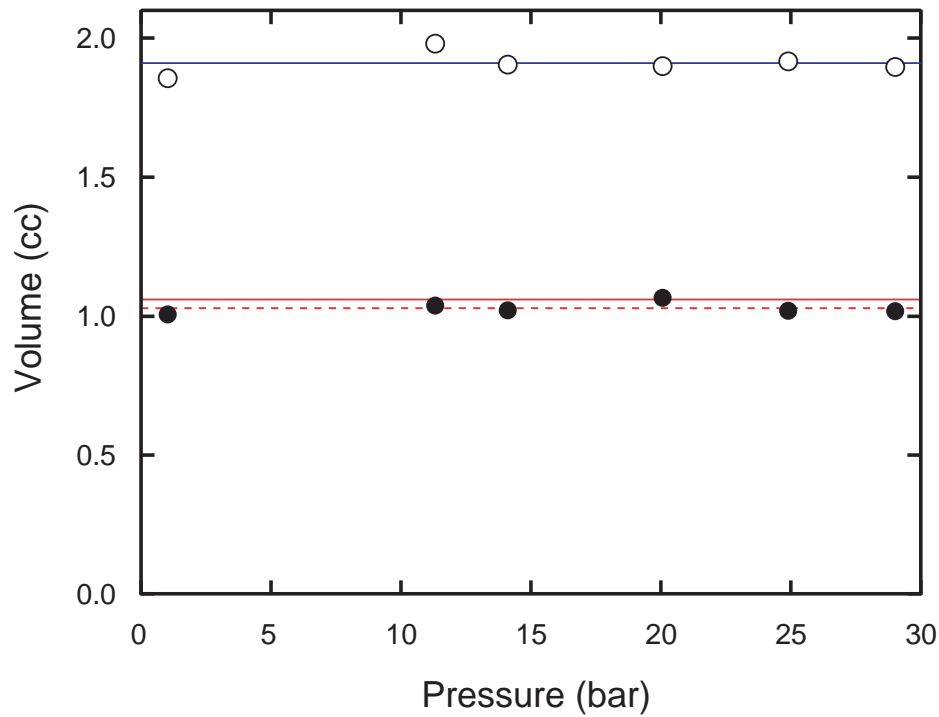
- Phase diagram for 98% aerogel/superfluid ^3He is the B phase at zero field at all pressures except close to T_C .
- The PCP has vanished for 98% aerogel.
- The Inhomogeneous Isotropic Scattering Model is consistent with the B-phase in aerogel, effect of magnetic field and heat capacity with $\lambda \approx 150$ nm; $\xi_a \approx 40$ nm.
- Heat capacity indicates that quenched disorder from aerogel produces gapless superfluidity (Andreev bound states).
- Scattering anisotropy may stabilize the equal-spin pairing state, ie A-phase, and proposal for stabilizing this phase macroscopically by uniaxial strain.

*Aerogel/Superfluid
 ^3He Heat Capacity:*

Calorimeter



Addendum from paramagnetic adsorbed solid ^3He

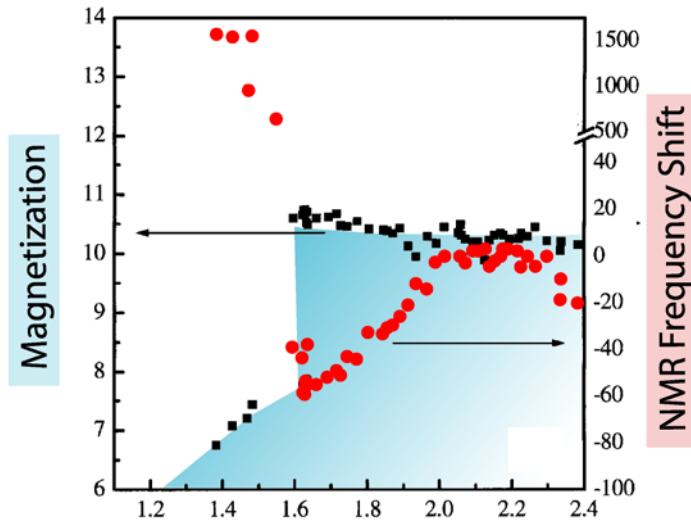


Geometric volume of the aerogel sample compared to the measured volume are within 3%.

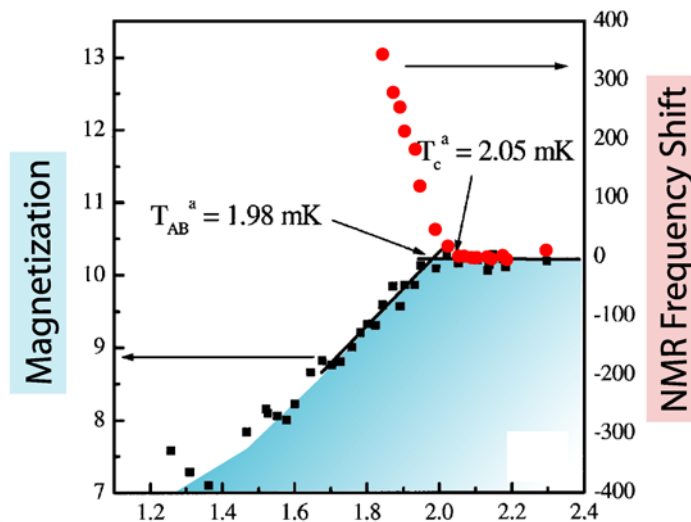
Is there an AB transition for superfluid ^3He in aerogel ?

First observed on supercooling by the **Stanford group** *Barker et al. PRL 85, 2148 (2000)*, from NMR magnetization and frequency shift.

Cooling
P = 32 bar
B = 284 G



Warming
P = 32 bar
B = 284 G



Similar NMR results from,

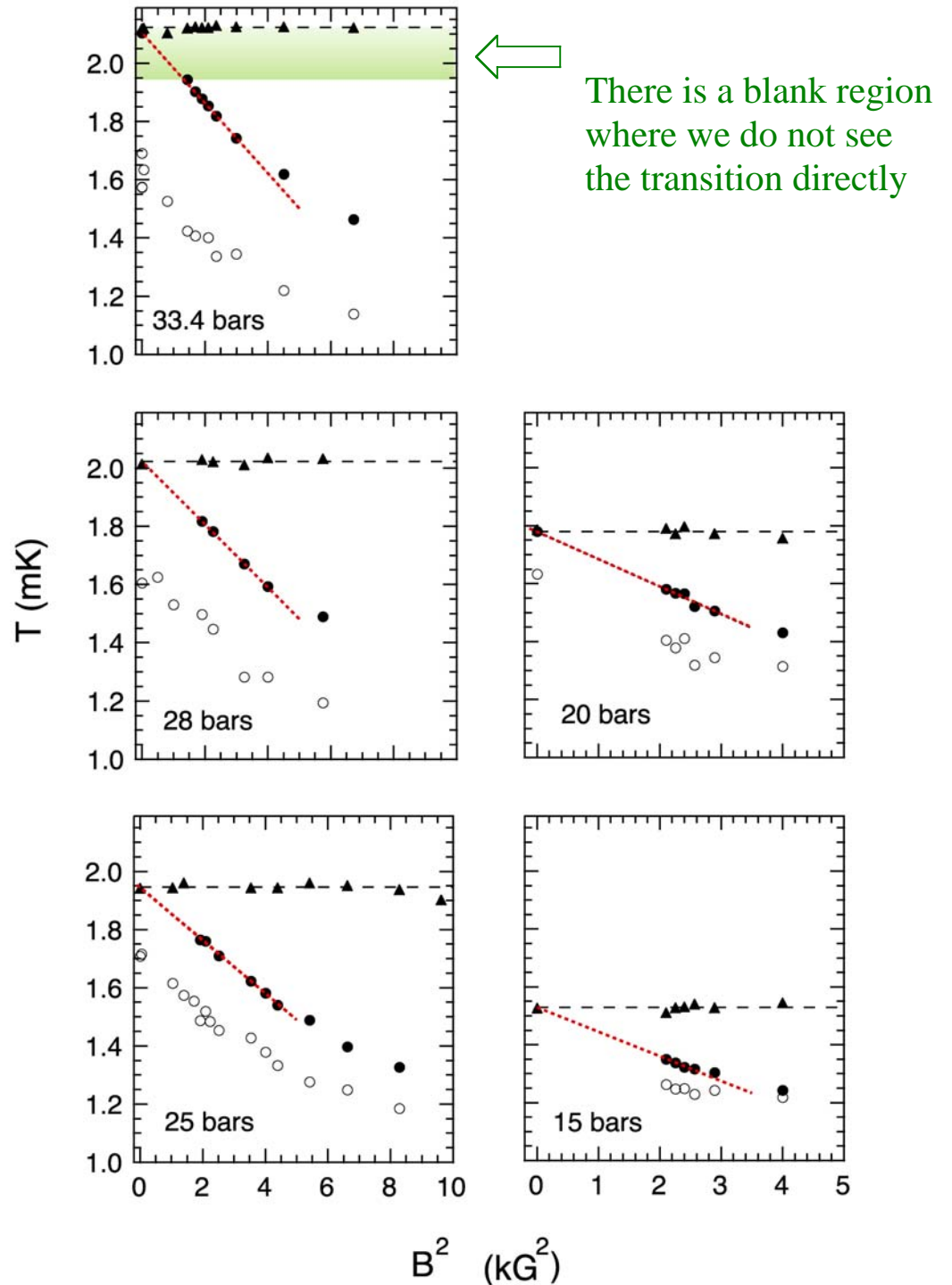
Northwestern, *Sprague et al. PRL 77, 4568 (1996)*;

and Grenoble (*E. Collin, PhD thesis, 2002*)

Aerogel superfluid ^3He :

AB transition was first observed by *Barker et al. PRL 85, 2148 (2000)* from NMR susceptibility and frequency shift

Transverse acoustic impedance, Gervais et al. PRL 87, 035701 (2001); PRB 66, 054528 (2002)



Metastable Phase Mixtures

Detail of how the tracking method allows exploration of the phase mixture near T_C

*Gervais et al., PRL 87, 035701 (2001);
PRB 66, 054528 (2002);*

